Chromium Substituted LiSn$_2$P$_3$O$_{12}$ Solid Electrolyte

R. Norhaniza$^1$, R.H.Y. Subban$^2$, N.S. Mohamed$^{3,*}$, A.Ahmad$^{4,*}$

$^1$ Institute of Graduate Studies, University of Malaya, 50603 Kuala Lumpur, Malaysia.
$^2$ Faculty of Applied Sciences, Universiti Teknologi MARA, 40450 Selangor, Malaysia
$^3$ Center for Foundation Studies in Science, University of Malaya, 50603 Kuala Lumpur, Malaysia.
$^4$ School of Chemical Science and Food Technology, Faculty of Science and Technology, Universiti Kebangsaan Malaysia 43600 Bangi Selangor Malaysia.
*E-mail: nsabirin@um.edu.my, azizan@ukm.my

Received: 25 June 2012 / Accepted: 2 September 2012 / Published: 1 October 2012

This study was focused on enhancing the conductivity of lithium tin phosphate, LiSn$_2$P$_3$O$_{12}$ by partially substituting Sn$^{4+}$ ions with Cr$^{3+}$ ions to obtain system with general formula Li$_{1+x}$Cr$_x$Sn$_{2-x}$P$_3$O$_{12}$ powders with $x = 0.2$, 0.4, 0.6 and 0.8 were prepared by mechanochemical milling method. X-ray diffraction analysis indicated that all samples consisted of triclinic crystalline LiSn$_2$P$_3$O$_{12}$ structure. Energy Dispersive X-ray analysis suggested that Cr$^{3+}$ was successfully substituted into the LiSn$_2$P$_3$O$_{12}$ crystalline structure. Impedance analysis showed an increase in conductivity with increase in $x$. The enhancement in bulk conductivity was due to increase in number and polarisability of Li$^+$ ions and crystallinity. The increase in grain boundary conductivity was ascribed to enhancement in Li$^+$ migration as a result of increase in contact between grains. The conductivity versus the reciprocal of temperature plots showed a sudden change in conductivity at 393 K for the sample with $x = 0.2$ and 373 K for the sample with $x = 0.4$, 0.6 and 0.8. This was attributed to a structure transition phenomenon. All samples showed Arrhenian behaviour.

Keywords: Tin, Conductivity, Nasicon, Arrhenius, Nanomaterial, Solid Electrolyte

1. INTRODUCTION

Extensive studies on lithium ion conductor with NASICON (Sodium Superionic Conductor) structure have been done because of their possible application in electrochemical devices [1-4]. NASICON has chemical formula LiM$_2$P$_3$O$_{12}$ ($M = $ Ti, Ge, Sn, Zr, Hf, etc.) with structure built from MO$_6$ octahedra and PO$_4$ tetrahedra. Their framework is formed by the corner sharing of oxygen of PO$_4$ tetrahedra with two MO$_6$ octahedra. In this framework there are two probable sites for Li$^+$ ions to
occupy i.e.: (i) M$_1$ sites located between MO$_6$ octahedral and (ii) M$_2$ sites, located between the ribbons of PO$_4$ tetrahedra. This 3D structure possesses channels which are also known as bottlenecks that allow migration of mobile ions. The corner sharing bonding in this framework also allows rotation of the framework; hence modifications like partial substitution may be done to improve its properties [2-4].

In the NASICON family, LiSn$_2$P$_3$O$_{12}$ has received less attention since their pellets were fragile and easily broken making them difficult to handle during characterisation [4-7]. This was due to structure transition (monoclinic to rhombohedral structure) which yielded changes in the lattice volume leading to breakage of the pellet. Several works have been done to overcome this problem such as the use of teflon as a binder [6-7] or the use of germanium (Ge) as an agent for completing the reaction to produce LiSn$_2$P$_3$O$_{12}$ [8]. In our previous study, we have proven that mechanically stable LiSn$_2$P$_3$O$_{12}$ pellets could be obtained by preparing this compound using mechanochemical milling technique without the use of any binding agent [9]. However, the LiSn$_2$P$_3$O$_{12}$ pellets showed low conductivities due chemical composition and ionic mobility.

The present study focussed on enhancing the conductivity of LiSn$_2$P$_3$O$_{12}$ by partially substituting Sn$^{4+}$ ions with Cr$^{3+}$ ions. The effects of partial substitution using Cr$^{3+}$ on conductivity of NASICON structured compounds have been investigated by several researchers [10-13]. According to Zhao et al [14], partial substitution can be easily carried out if the size of the ion is about the same as that of the ion to be substituted. Since the ionic radius of Cr$^{3+}$ ion is 0.62 Å while that of Sn$^{4+}$ ion is 0.65 Å, therefore partial substitution of Sn$^{4+}$ by Cr$^{3+}$ is expected to be possible. Xu et al reported on conductivity enhancement of about three orders of magnitude for Cr$^{3+}$ substituted LiGe$_2$P$_3$O$_{12}$ system [10]. Delmas et al [11] reported that the conductivity enhancement in Cr$^{3+}$ substituted NaZrP$_3$O$_{12}$ was due to increase in mobility of ions as reflected by decrease in activation energy. This partial substitutions were done with the objectives of creating Li interstitials to form structures with general formula Li$_{1+x}$Cr$_x$Sn$_{2-x}$P$_3$O$_{12}$ where Sn$^{4+}$ ↔ Li$^+$ + Cr$^{3+}$ [10-14]. The effects of Cr substitution on the structure and conductivity of LiSn$_2$P$_3$O$_{12}$ are reported in this article.

2. EXPERIMENTAL

2.1. Materials and Methods

Stoichiometric proportions of Li$_2$O (99 % Aldrich), Cr$_2$O$_3$, SnO$_2$ (98 % Sigma-Aldrich) and NH$_4$H$_2$PO$_4$ (98 % Sigma-Aldrich) were mixed according to the formula Li$_{1+x}$Cr$_x$Sn$_{2-x}$P$_3$O$_{12}$ ($x$ = 0.2, 0.4, 0.6 and 0.8) and heated at 700 °C for two hours. The heating process was required in order to decompose hydrogen phosphate and initiate reaction with carbonate [13,15]. After the heating process, the mixtures were milled for 80 hours using Fristch 7 planetary ball mill. The milled powders were then pelletized to form pellets with diameter of 13 mm and thickness of ~ 1 mm using a hydraulic press at a pressure of 7 ton. The pellets were then sintered at four different temperatures from 700 to 1000 °C for 8 hours. However, it was observed that for different $x$ values, optimum conductivity was achieved at different sintering temperatures which were at 700, 800, 900 and 1000 °C for $x$ = 0.2, 0.4, 0.6 and 0.8 respectively. Therefore, for further analyses, samples were prepared by sintering them at
their optimum sintering temperatures and these samples are labelled as C2, C4, C6 and C8 for \( x = 0.2, 0.4, 0.6 \) and 0.8 respectively.

Structural study was carried out using Bruker D8 Advance X-Ray Diffractometer. Crystallite size, \( D \) was calculated using Scherer equation:

\[
D = \frac{k\lambda}{\beta \cos \theta}
\]

where \( k \) is the Scherer constant value (0.94), \( \lambda \) is the wavelength of the source (Copper-K, \( \lambda = 1.5406\text{Å} \)), \( \beta \) is FWHM (in radians) and \( \theta \) is the Bragg angle (in radians). Surface morphology and elemental composition analyses were done using FEI Quanta 200 FESEM equipment in low vacuum condition and accelerating voltage of 20 kV. For these analyses, samples were gold sputtered after ultrasound cleaning. Conductivity studies were carried out using Solartron 1260 impedance analyser. Transference number was determined using Wagner’s polarisation method on Solartron 1260 interfaced with SI 1287 Electrochemical Interface. The samples were sandwiched between two carbon electrodes and polarised by applying a potential of 0.5 V.

3. RESULTS AND DISCUSSION

3.1. X-ray diffraction (XRD) study

Fig. 1 shows XRD spectra of samples C2, C4, C6 and C8. Presence of crystalline triclinic LiSn\(_2\)P\(_3\)O\(_{12}\) (labelled with \( hkl \) index in Fig. 1) was detected in all samples. The values of lattice constants of this structure for all samples were similar with those reported by Iglesias which were \( a = 14.680 \text{Å}, b = 8.413 \text{Å} \) and \( c = 8.900 \text{Å} \) [16]. Our studies [9] on unsubstituted LiSn\(_2\)P\(_3\)O\(_{12}\) showed presence of rhombohedral structure with lattice constant values of \( a = 8.650 \text{Å}, b = 8.650 \text{Å} \) and \( c = \)
21.487 Å. The change from rhombohedral structure in the unsubstituted LiSn₂P₃O₁₂ to triclinic structure in the chromium substituted samples may be due to the presence of interstitial Li⁺ ions at the mid way sites of M₁ and M₂ [1,3]. The change in structure may also be attributed to the restructuring of the lattice atoms in order to accommodate Cr³⁺ ions which were smaller compared to Sn⁴⁺ ions. The adjustment in lattice parameters was due to the expansion of the framework along the a-axis which caused a rotation of tetrahedral PO₄ units connecting two parallel columns with a concomitant contraction of the structure along c-axis [17,18].

Figure 2. Variation of crystallite size with x (chromium content) at different planes

The peaks of the LiSn₂P₃O₁₂ triclinic crystalline structure were also observed to become narrower with increase in Cr³⁺ content showing that Cr substitution controlled the number of nucleation centres resulting in the formation of larger crystallites [19]. Increase in XRD intensity with x value was due to the growth of LiSn₂P₃O₁₂ structure [10, 20]. In order to determine the crystallite size, Scherer equation was used and the values of crystallite sizes for LiSn₂P₃O₁₂ are shown in Fig. 2. C₈ possessed the largest crystallite size in the range 37 to 38 nm while C₂, C₄ and C₆ possessed crystallites in the range from 11 to 30 nm, 12 to 33 nm and 14 to 34 nm, respectively. This inferred that samples were nanomaterials. Fig. 2 also shows that the crystallite size in C₈ was homogeneous while those in other samples were quite heterogeneous. Traces of SnP₂O₇ crystalline indicated by the peak at 2θ ~ 23 ° was found to be present in all samples. However, the amount of SnP₂O₇ decreased with increase in x. Meanwhile, trace of Li₆P₆O₁₈ was also found to be present in C₄, C₆ and C₈ as indicated by the presence of the peak at 2θ ~ 17.5 °.

3.2. Scanning Electron Microscopy (SEM)

Fig. 3 illustrates the SEM micrograph for the cross-sections of C₂, C₄, C₆ and C₈. Average grains size in the samples increased with Cr content which were 0.12 μm, 0.45 μm, 0.49 μm and 0.67 μm for C₂, C₄, C₆ and C₈ respectively. A notable difference in the grain size was observed in C₂.
Decrease in number of large grains was seen in C4. Increase in number of grains with similar size and also contacts between grains were noticed in C6 while C8 exhibited very dense and well-packed structure with fewer pores. Uniform grains that fused together were also observed in this sample.

Figure 3. SEM images of Li$_{1+x}$Cr$_x$Sn$_{2-x}$P$_3$O$_{12}$ samples

3.3. Energy Dispersive X-ray Spectroscopy (EDX)

EDX analysis was done at different points (a and b indicated in Fig. 3) to confirm stoichiometric proportions of the samples and the EDX backscattered spectra are shown in Fig. 4. Table 1 lists the atomic percentage at both points a and b in Fig. 3. Since Li$^+$ ion was not detectable by EDX, the concept of charge neutrality was employed [20-21]. Using this concept, it was found that the stoichiometries of the samples were in good agreement with the designated values. Thus EDX analysis indicated that Cr$^{3+}$ was successfully substituted into LiSn$_2$P$_3$O$_{12}$ matrix forming Li$_{1+x}$Cr$_x$Sn$_{2-x}$P$_3$O$_{12}$ compounds.
Figure 4. EDX spectra at both points a and b for all samples.
Table 1. Atomic percentage of elements at points a and b for all samples

<table>
<thead>
<tr>
<th>Sample</th>
<th>Composition</th>
<th>Elements (Atomic %)</th>
<th>Stoichiometric ratio Cr : Sn : P</th>
</tr>
</thead>
<tbody>
<tr>
<td></td>
<td></td>
<td>Cr</td>
<td>Sn</td>
</tr>
<tr>
<td>C2</td>
<td>starting mixture</td>
<td>4</td>
<td>36</td>
</tr>
<tr>
<td></td>
<td>EDX analysis C2-a</td>
<td>3.56</td>
<td>35.54</td>
</tr>
<tr>
<td></td>
<td></td>
<td>4.15</td>
<td>36.35</td>
</tr>
<tr>
<td>C4</td>
<td>starting mixture</td>
<td>7.98</td>
<td>32.02</td>
</tr>
<tr>
<td></td>
<td>EDX analysis C4-a</td>
<td>7.26</td>
<td>32.20</td>
</tr>
<tr>
<td></td>
<td></td>
<td>7.55</td>
<td>32.17</td>
</tr>
<tr>
<td>C6</td>
<td>starting mixture</td>
<td>12.01</td>
<td>28.03</td>
</tr>
<tr>
<td></td>
<td>EDX analysis C6-a</td>
<td>11.46</td>
<td>28.26</td>
</tr>
<tr>
<td></td>
<td></td>
<td>11.94</td>
<td>27.92</td>
</tr>
<tr>
<td>C8</td>
<td>starting mixture</td>
<td>16.00</td>
<td>24.00</td>
</tr>
<tr>
<td></td>
<td>EDX analysis C8-a</td>
<td>15.45</td>
<td>23.66</td>
</tr>
<tr>
<td></td>
<td></td>
<td>15.86</td>
<td>23.71</td>
</tr>
</tbody>
</table>

3.4. Bulk Conductivity

Impedance plot (Z plot) for the sample with $x = 0.2$ is presented in Fig. 5. The plot shows the presence of overlapping semicircles and an arc. The values of bulk resistance ($R_b$) was determined from the intercept of the high frequency semicircle with the $Z'$-axis. The bulk ($\sigma_b$) conductivity was calculated using the obtained $R_b$ value and they are listed in Table 2. Since the semicircles in the impedance plot were quite difficult to be distinguished, confirmation of the obtained bulk conductivity was done by plotting the conductivity spectrum, $\sigma$ plot (imaginary conductivity, $\sigma''$ versus real conductivity, $\sigma'$) of the sample. This method has been used by a few groups of researchers [22-26].
The $\sigma''$ versus $\sigma'$ for C2 is depicted in Fig. 6. The figure shows the existence of two linear sections with different gradients (at high frequency region) and a semicircle (at low frequency region). Extrapolation of the first steeper section to $\sigma'$ axis gave the value of bulk conductivity. Similar method was used to obtain conductivity for other samples.

It was observed that the values of bulk conductivity obtained from both impedance and conductivity spectra, Table 2, were marginally close to each other hence confirming the values of $\sigma_b$ obtained from the impedance plots. Fig. 7 shows that both bulk and grain boundary conductivities
increased with increase in chromium content. The highest $\sigma_b$ and grain boundary conductivity, $\sigma_{gb}$ were observed for C8 with values of $5.62 \times 10^{-6}$ S cm$^{-1}$ and $1.44 \times 10^{-6}$ S cm$^{-1}$ respectively. This means that C8 showed an increase of three times of $\sigma_b$ and about one order of magnitude of $\sigma_{gb}$ of the unsubstituted LiSn$_2$P$_3$O$_{12}$ which were $1.65 \times 10^{-6}$ and $1.05 \times 10^{-7}$ S cm$^{-1}$, respectively [9].

The increase in conductivity with chromium content may be due to increase in the number of interstitial Li$^+$ ions due to increase in the amount of Cr$^{3+}$ ions [13]. Another factor that contributed to bulk conductivity enhancement was the increase in crystallinity which was the conducting structure in the compounds [13,19].

According to Semanate et al [20], the higher the structural order, the lower the energy barriers for ions to overcome in order to move. The enhancement in grain boundary conductivity was also attributable to improvement in contact between grains due to increase in homogeneity of grains as discussed in section 3.2. Semanate et al [20] also reported that similar grain sizes resulted in increase in the interfacial areas between grains. This increased the contact between grains which in turn increased conduction path and hence enhancing both bulk and grain boundary conductivities. Similar observations have reported by Best et al [2,27] and Mei et al [28].

The variations of conductivity with temperature for all samples are depicted in Fig. 8. For C2, $\sigma_b$ showed two linear regions with different slopes with a sudden change in slope at 393 K. The linear regions showed that temperature dependent conductivity was thermally activated indicating Arrhenius behaviour [29-31]. The same trend was observed for other samples. However, the sudden change in slope was observed at 373 K. The bulk activation energy, $E_{a,b}$ values are tabulated in Table 2. All samples showed similar $E_{a,b}$ values at low temperature (low T) and high temperature (high T) regions. However the $E_{a,b}$ in high T was lower than low T suggesting that there were two different diffusion channel [20,31].

![Figure 8. Temperature dependence of conductivity for Li$_{1+x}$Cr$_x$Sn$_{2-x}$P$_3$O$_{12}$ samples](image-url)
This was associated with a structure transition phenomenon from triclinic to rhombohedral structure [32]. Such structural transition occurred due to redistribution of Li$^+$ ions. Losilla and co-workers [4] described that low temperature triclinic structure had low symmetry with ordered Li$^+$ ions distribution. The structure had several bottleneck positions with a narrow probability for the Li$^+$ ions to conduct. As temperature increased, Li$^+$ bonds became loose making it free to move and became randomly distributed. At the same time, the torsion on PO$_4$ tetrahedra takes place due to the frameworks accommodating themselves without breaking their bonds [17]. As a result PO$_4$ tetrahedra rotated around two fold symmetry while SnO$_6$ octahedra rotated in the opposite direction to its direction of rotation at low temperature forming restructured frameworks [33]. These restructured frameworks possessed a more appropriate bottleneck arrangement for the hopping of Li$^+$ ions to take place. According to Kolher and Schulz [34], with increasing temperature, the bottleneck size increased hence ideally for diffusion of Li$^+$ ion. As such, the probability of the Li$^+$ ions to move increased causing redistribution of Li$^+$ ions into lattice sites. This redistribution of lithium ions gave rise to a formation of a new structure (rhombohedral) which was a more conducting structure.

Table 2. Comparison of bulk conductivity values of Li$_{1+x}$Cr$_x$Sn$_{2-x}$P$_3$O$_{12}$ samples obtained from impedance and conductivity plots together with activation energy and transference number values.

<table>
<thead>
<tr>
<th>x</th>
<th>$\sigma_b \times 10^{-6}$ (S cm$^{-1}$)</th>
<th>$E_{a,b}$ (±0.01) eV</th>
<th>$\tau_{ion}$</th>
</tr>
</thead>
<tbody>
<tr>
<td></td>
<td>Z plot</td>
<td>$\sigma$ plot</td>
<td>Low T</td>
</tr>
<tr>
<td>0.2</td>
<td>2.13 (± 0.03)</td>
<td>2.11 (± 0.01)</td>
<td>0.10</td>
</tr>
<tr>
<td>0.4</td>
<td>2.34 (± 0.02)</td>
<td>2.35 (± 0.04)</td>
<td>0.09</td>
</tr>
<tr>
<td>0.6</td>
<td>3.97 (± 0.02)</td>
<td>3.96 (± 0.03)</td>
<td>0.09</td>
</tr>
<tr>
<td>0.8</td>
<td>5.62 (± 0.02)</td>
<td>5.61 (± 0.01)</td>
<td>0.07</td>
</tr>
</tbody>
</table>

The values of $E_{a,b}$ also indicated that mobility of ions was not varied. As such, the increase in conductivity may be due to increase in polarisability of Cr ions and possible disorder of occupancy of vacant sites. Similar observation has been reported by Chowdari et al [29]. While Sugantha and Varadaraju [30] reported that the larger the polarisability of cation-oxygen bond, the easier the transport of ions through the bottleneck. Therefore it can be summarised that increase in Cr$^{3+}$ ions, increased the ion transportation which resulted in an increase in conductivity.

3.5. Transference Number Measurement

Fig. 9 shows a typical plot of normalized polarisation current versus time for Li$_{1+x}$Cr$_x$Sn$_{2-x}$P$_3$O$_{12}$ system. The polarisation current for all samples saturated after about 50 s as tabulated in Table 2. All samples showed transference number value close to unity. This shows that the conductivity of the Li$_{1+x}$Cr$_x$Sn$_{2-x}$P$_3$O$_{12}$ electrolytes was predominantly ionic and was expected to be that of Li$^+$ ions.
4. CONCLUSIONS

Structural studies showed increase in crystallinity and homogeneity of the Cr substituted LiSn$_2$P$_3$O$_{12}$ with increase in Cr content. EDX analysis indicated that Cr was successfully substituted at Sn sites in LiSn$_2$P$_3$O$_{12}$ matrix. This was supported by the XRD results that showed changes in lattice constants as compared to the unsubstituted LiSn$_2$P$_3$O$_{12}$. SEM results showed increase in contact between grains due to grain with similar size when Cr is substituted into the parent compound. The substitution also resulted in conductivity enhancement up to to the order of $10^{-6}$ S cm$^{-1}$. Temperature dependent conductivity studies showed change in crystallography structure as reflected by different $E_{ab}$ values for high and low temperatures ranges. The enhancement in bulk conductivity with Cr content was attributed to increase in crystallinity and homogeneity of LiSn$_2$P$_3$O$_{12}$ structure. Other factors of this enhancement were increase in number and polarisability of Li$^+$ ions. The enhancement in both conductivities were also due to increase in Li$^+$ migration as a result of increase in contact between grains.

References


© 2012 by ESG (www.electrochemsci.org)